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# 01. Introduction: Maps

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## Abstract

Part one of course materials for Nonequilibrium Statistical Physics (Physics 626), taught by Gerhard Müller at the University of Rhode Island. Entries listed in the table of contents, but not shown in the document, exist only in handwritten form. Documents will be updated periodically as more entries become presentable.

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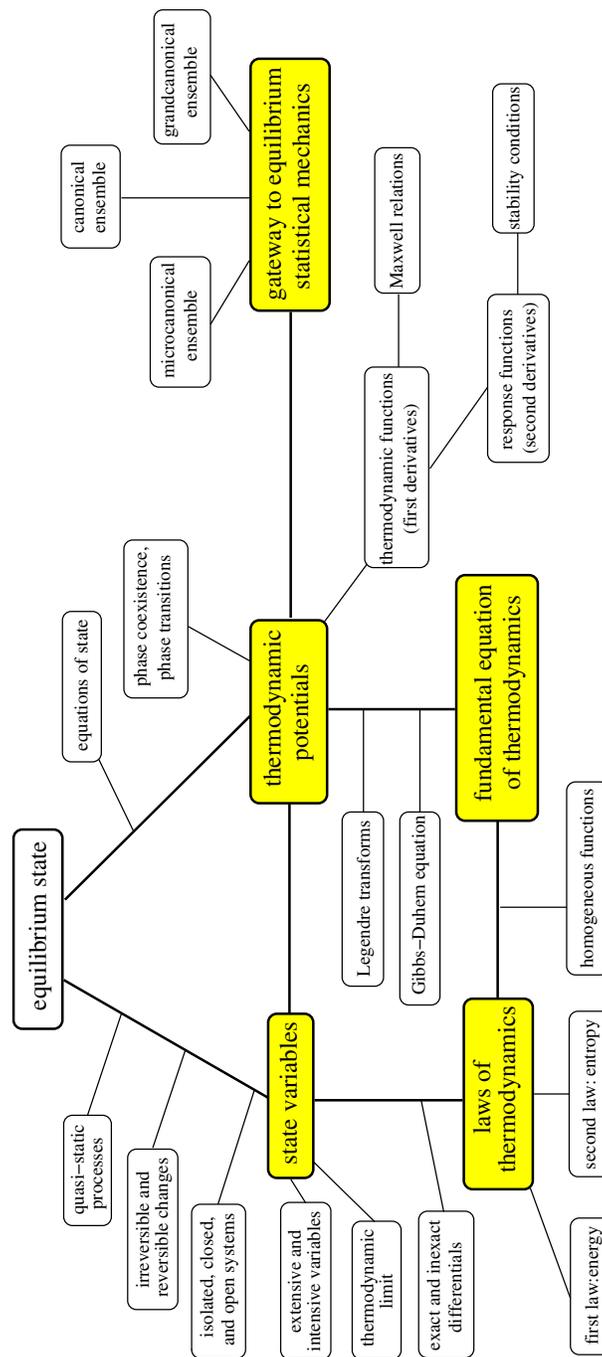
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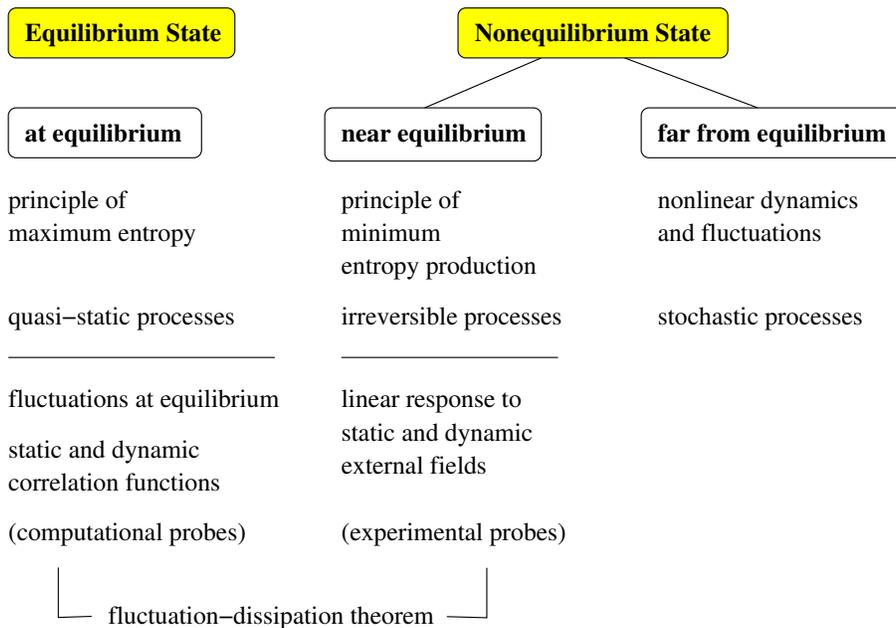
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# Equilibrium Thermodynamics Overview [nlm6]



# Thermal Equilibrium and Nonequilibrium [nlh1]

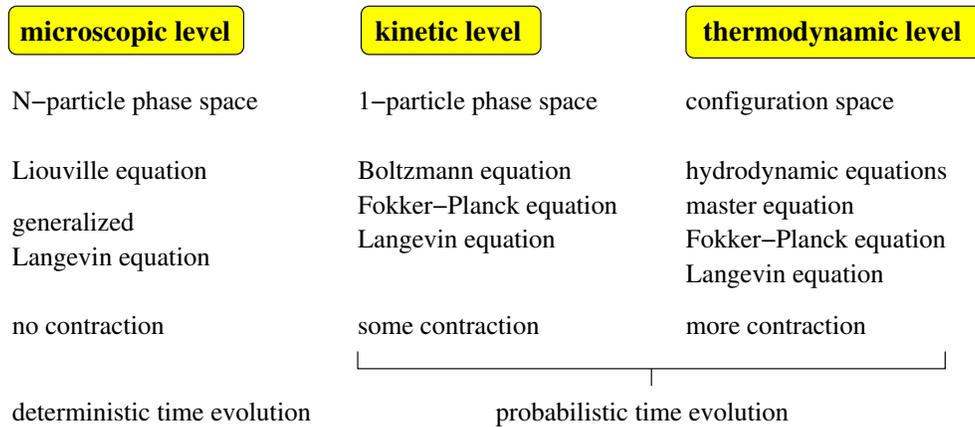


Distinguish independently between

- *equilibrium* and *nonequilibrium* situations,
- *time-independent* and *time-dependent* phenomena.

	<b>equilibrium situation</b>	<b>nonequilibrium situation</b>
<b>time-independent phenomena</b>	equal-time correlations	equal-time correlations in steady states
<b>time-dependent phenomena</b>	delayed-time correlations	delayed-time correlations in steady states any correlations in non-steady states

# Levels of Description in Statistical Physics [nln2]



# Contraction – memory – time scales [nlh15]

microscopic dynamics	$\Rightarrow$ contraction $\Rightarrow$	stochastic dynamics
future state determined by present state alone	focus on subset of dynamical variables	future state determined by present and past states
deterministic time evolution of dynamic variables	$\Downarrow$	ignoring memory of past makes dynamics of selected variables probabilistic
	judicious choice: slow variables and long time scales	deterministic time evolution of probability distributions and mean values
	$\Downarrow$ $\Rightarrow$	short memory of fast variables has little impact on dynamics of slow variables at long times

## Degrees of memory:

- **Future independent of present and past.**

Completely factorizing process.

$$P_j(x_1, t_1; x_2, t_2; \dots) = P_s(x_1, t_1) P_s(x_1, t_1) \dots$$

- **Future dependent on present only.**

Markov process.

$$P_c(x_1, t_1; x_2, t_2; \dots | \bar{x}_1, \bar{t}_1; \bar{x}_2, \bar{t}_2, \dots) = P_c(x_1, t_1; x_2, t_2; \dots | \bar{x}_1, \bar{t}_1)$$

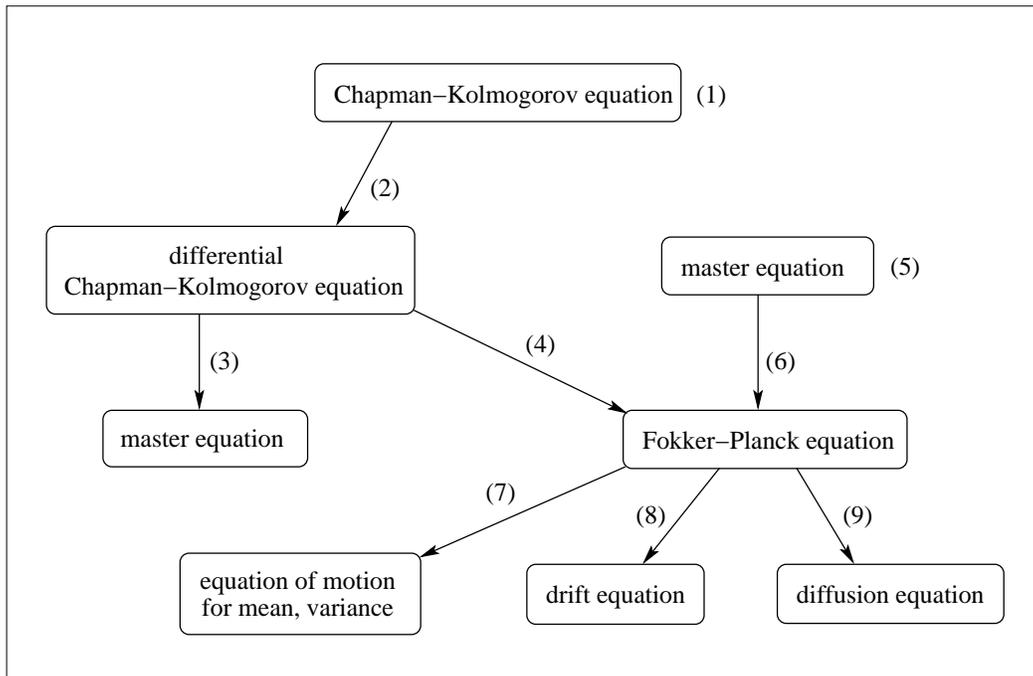
for  $t_1 \geq t_2 \geq \dots \geq \bar{t}_1 \geq \bar{t}_2 \geq \dots$ .

- **Future dependent on present and past.**

Non-Markovian process.

$$P_c(x_1, t_1; x_2, t_2; \dots | \bar{x}_1, \bar{t}_1; \bar{x}_2, \bar{t}_2, \dots).$$

# Markov processes: map of specifications [mln16]



- (1) Chapman-Kolmogorov equation imposes restrictions on permissible functions  $P(x, t|x_0)$  but does not suggest a classification of processes.
- (2) Particular solutions that are specified by
  - $A(x, t)$  describing drift,
  - $B(x, t)$  describing diffusion,
  - $W(x|x'; t)$  describing jumps.
- (3) Jump processes exclusively.
- (4) Processes with continuous sample paths, satisfying Lindeberg criterion (drift and diffusion, no jumps).
- (5) Master equation with any  $W(x|x'; t)$  specifies a Markov process. Natural starting point for processes with discrete stochastic variables.
- (6) Transition rates  $W(x|x'; t)$  of master equation approximated by two jump moments provided they exist. Approximation captures drift and diffusion parts (on some scale).
- (7) Drift and diffusion determine mean  $\langle\langle x \rangle\rangle$  and variance  $\langle\langle x^2 \rangle\rangle$  via equations of motion for jump moments.
- (8) Deterministic process have no diffusive part:  $B(x, t) = 0$ .
- (9) Purely diffusive processes have no drift:  $A(x, t) = 0$ .

# Brownian motion: panoramic view [nln23]

- Levels of contraction (horizontal)
- Modes of description (vertical)

→ contraction →

relevant space	$N$ -particle phase space	1-particle phase space	configuration space
dynamical variables	$\{\mathbf{x}_i, \mathbf{p}_i\}$	$\mathbf{x}, \mathbf{p}$	$\mathbf{x}$
theoretical framework	Hamiltonian mechanics	Langevin theory	Einstein theory
... for dynamical variables	generalized Langevin equation	Langevin equation (for $dt \ll \tau_R$ )	Langevin equation (for $dt \gg \tau_R$ )
... for probability distribution	quant./class. Liouville equation	Fokker-Planck equation (Ornstein-Uhlenbeck process)	Fokker-Planck equation (diffusion process)

- Here  $dt$  is the time step used in the theory and  $\tau_R$  is the relaxation time associated with the drag force the Brownian particle experiences.
- The generalized Langevin equation is equivalent to the Hamiltonian equation of motion for a generic classical many-body system and equivalent to the Heisenberg equation of motion for a generic quantum many-body system.

# Linear response and equilibrium dynamics [nlh24]

