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Post Eruption inflation of the East Pacific Rise at 9° 50′ N

Scott L. Nooner1, Spahr C. Webb2, W. Roger Buck2, and Marie-Helene Cormier3

1Department of Geography and Geology, University of North Carolina Wilmington, Wilmington, North Carolina, 2Lamont-Doherty Earth Observatory, Columbia University, Palisades, New York, 3University of Rhode Island, Graduate School of Oceanography, Narragansett, Rhode Island

Abstract In June 2008, we installed a geodetic network at 9° 50′ N on the East Pacific Rise to track the long-term movement of magma following the 2005/6 eruption. This network consists of 10 concrete benchmarks stretching from the ridge to 9 km off-axis. During three campaign-style surveys, measurements of vertical seafloor motions were made at each of these benchmarks by precisely recording ambient seawater pressure as a proxy for seafloor depth with a mobile pressure recorder (MPR). The MPR was deployed using the manned submersible Alvin in 2008 and 2009 and the remotely operated vehicle Jason in 2011. The MPR observations are supplemented with data from a multiyear deployment of continuously recording bottom pressure recorders (BPRs) extending along this segment of the ridge that can record rapid changes in seafloor depth from seafloor eruptions and/or dike intrusions. These measurements show no dike events and up to 12 cm of volcanic inflation that occurred from December 2009 to October 2011 in the area of the 2005/6 eruption. These observations are fit with an inflating point source at a depth of 2.7 km and volume change of $2.3 \times 10^6 \text{m}^3/\text{yr}$ located on the ridge axis at approximately 9° 51.166′ N, 407 m from our northernmost benchmark, suggesting that the magma chamber underlying this segment of the ridge is being recharged from a deeper source at a rate that is about half the long-term inflation rate observed at Axial Seamount on the Juan de Fuca Ridge. These data represent the second location that active volcanic uplift has been measured on a mid-ocean ridge segment, and the first on a nonhotspot influenced segment.

1. Introduction

The East Pacific Rise (EPR) at 9° 50′ N is a fast spreading mid-ocean ridge (MOR) separating the Pacific and Cocos plates (Figure 1) which move away from each other at a full rate of about 110 mm/yr [Carbotte and McDonald, 1992] while the entire ridge system migrates to the north-west at 50 mm/yr relative to a hotspot reference frame [Gripp and Gordon, 2002]. It is one of the best-studied MOR segments in the world and it supports a dynamic hydrothermal system fueled by repeated dike intrusions and eruptions. The first documented eruption occurred in 1991/1992 [Haymon et al., 1993; Rubin et al., 1994], and since that time scientists have regularly returned to document changes in ecosystem, vent-fluid chemistry, and temperature, and to conduct detailed geological mapping. Because of these studies, when another eruption occurred in 2005/2006 [e.g., Fornari et al., 2012; Tolstoy et al., 2006], seismicity changes associated with the event were recorded on a local array of Ocean Bottom Seismometers (OBSs). Using a submersible, Soule et al. [2007] mapped out the extent of new lava flows in a 17 km long region of the EPR (Figure 1). These eruptive events drive changes in the composition of vent fluids [Fornari et al., 2012; Von Damm, 2004, p. 2013; Yücel and Luther, 2013] and affect the temporal succession of biological communities surrounding the vents [Luther et al., 2012; Shank et al., 1998]. Between these diking events, magma is presumably redistributed subsurface and the underlying magma lens may be refilled, causing changes in depth of the seafloor above, a behavior that has previously been measured only at Axial Seamount on the Juan de Fuca Ridge [Chadwick et al., 2012, 2006; Nooner and Chadwick, 2009]. The presence of a shallow axial magma body a few tens of meters thick and 0.5-4.5 km wide within the EPR was first established from a multichannel seismic (MCS) experiment in 1985 [Detrick et al., 1987; Kent et al., 1993a, 1993b; Vera et al., 1990]. 3-D MCS data acquired in 2008 extend the earlier observations and have allowed fine-scale details regarding crustal magma distribution to be revealed [Carbotte et al., 2012]. Magma lens segments 5–20 km in length are present along axis, with many segment ends corresponding to discontinuities in seafloor structures [Carbotte et al., 2013].
In this paper, we present results from a series of three geodetic surveys that were carried out at 9° 50′N on the EPR using water pressure measurements on top of a network of 10 concrete benchmarks that were emplaced on the seafloor in June 2008 (Figure 1). Our results allow us to compare magma recharge at the fast spreading EPR with the hot-spot influenced intermediate spreading rate Axial Seamount segment of the Juan de Fuca ridge—the only other submarine mid-ocean ridge site with documented magmatic uplift. We also compare results with observations made on recently-active, subaerial spreading centers in Iceland [Einarsson, 1991], and most recently the Afar of Ethiopia [e.g., Wright et al., 2012]. Post eruption deformation at these locations has been observed and interpreted as due to either magma recharge or viscoelastic relaxation [e.g., Nooner and Chadwick, 2009; Nooner et al., 2009].

The implementation of geodetic monitoring experiments in the marine environment has been a challenge. Seafloor geodetic techniques include acoustic ranging between pairs of instruments such as extensometers [Chadwell et al., 1999; Chadwick and Stapp, 2002; Chadwick et al., 1999; McGuire and Collins, 2013; Nagaya et al., 1999], combined GPS/acoustic positioning of instruments on the bottom from surface ships [e.g., Chadwell et al., 1995; Fujimoto et al., 2003; Fujita et al., 2003; Gagnon et al., 2005; Ito et al., 2011; Osada et al., 2003; Sato et al., 2011; Spiess et al., 1998], long and short baseline tiltmeters [Anderson et al., 1997; Tolstoy et al., 1998], and seafloor gravity measurements [e.g., Ballu et al., 1999, 2009, 1998; Nooner et al., 2007; Sassa-gawa et al., 2003]. Although each of these techniques has advantages, one of the simplest and most successful techniques has been the combined use of bottom pressure recorders (BPRs) and mobile pressure recorders (MPRs). BPRs sit on the seafloor and continuously record ambient pressure as a proxy for seafloor depth [Fox, 1999; Fujimoto et al., 2003; Watanabe et al., 2004] using sea level as a datum so that any uplift or subsidence of the seafloor causes a corresponding decrease or increase in measured pressure. The BPRs we used in this study were built at Lamont-Doherty Earth Observatory and use Paroscientific Digiquartz model.

Figure 1. Bathymetry map of the study area is shown with locations of seafloor benchmarks (gray circles), BPR deployment locations (squares), and the location of the best-fitting inflation source (dark blue star). Benchmarks EPR09, EPR07, and EPR10 are located just east of the ridge axis, less than 100 m from the Axial Summit Trough (AST). The outline from the 2005–2006 eruption [Soule et al., 2007] is indicated by the purple line.
410K pressure transducers. These gauges are subject to long-term linear drift in pressure of up to 150 ppm that can yield inferred height changes up to 38 cm/yr at a depth of 2500 meters, the depth of the axial summit trough (AST) [Chadwick et al., 2006; Fujimoto et al., 2003; Polster et al., 2009; Watts and Kontoyiannis, 1990]. The exact rate of drift is specific to each sensor and stabilizes during the first few weeks of deployment while the sensor equilibrates [Eble et al., 1989; Fox, 1990]. Because of the inherent drift, these instruments are very good for observing sudden, short period, or episodic events, but are inadequate for observing long-term or gradual deformation.

Campaign style measurement using MPRs [Chadwick et al., 2012, 2006; Nooner and Chadwick, 2009; Sasa-gawa et al., 2003; Stenvold et al., 2006] are analogous to optical leveling on land, but involve making pressure measurements on an array of fixed seafloor benchmarks. The MPR instrument consists of two Paroscientific Digiquartz pressure transducers (again, model 410K in this study). The advantage of this technique is that the drift in the pressure gauges can be calculated by assuming no relative deformation occurs during the survey (typically 2–3 days). The disadvantage of this technique is that all the measurements are made with respect to a reference benchmark, which is assumed to be stationary over time. This requires a site outside of the region of expected deformation. MPR uncertainties are typically less than 0.9 cm [Chadwick et al., 2012, 2006; Nooner and Chadwick, 2009] and have been as low at 0.5 cm [Nooner et al., 2007; Stenvold et al., 2006]. BPRs and MPR measurements are typically colocated to allow drift calibration of the BPRs.

Here we present results from BPR records and three MPR surveys at 9°50′N on the EPR, spanning 3.5 years, that show uplift occurring subsequent to the 2005–2006 eruption. We use the deformation to estimate the location of the inflation source and the magma supply rates and compare to what has been observed at Axial Seamount on the Juan de Fuca Ridge, Iceland and Ethiopia.

2. Methods and Results

A bottom pressure recorder (BPR08) was deployed near the 9°50′ N AST beginning in 2007. BPR08 was then recovered, serviced, and redeployed in both 2008 and 2009, colocated with MPR benchmarks EPR09 and EPR07, respectively. BPR08 stopped working shortly after being redeployed in 2009. Additional BPRs were deployed in 2009; BPR01 was deployed at benchmark EPR10 and BPR06 was deployed at benchmark EPR01 (Figure 1 and Table 1).

The BPRs collected data at a 40 Hz sample rate. The data were filtered to remove the variation in pressure caused by the ocean tides. Records were converted to depth using the standard ocean depth formula derived by Fofonoff and Millard [1983]. A detailed discussion of this procedure is given in Nooner [2005]. The resulting BPR records shown in Figure 2 are difficult to interpret because of gauge drift that varies from instrument to instrument and large (order 20cm) pressure signals of oceanographic origin. The oceanographic signals can be mostly removed by differencing adjacent sites (not shown), but the unpredictable gauge drift limits the usefulness of the BPR data for assessing slow geodetic changes. Instrument drift in BPR06 (order 1m) is clearly evident (Figure 2d) as it is unlikely to have experienced any uplift at all since it is colocated with the reference benchmark, EPR01, ~9 km from the AST. This illustrates the importance of combining BPR and MPR measurements in seafloor geodetic studies like this one.

Eruptions have been observed to produce large (order 1m) drops in seafloor elevation over time intervals of only a few hours as the underlying magma chamber deflates during the eruption [Chadwick et al., 2012, 1999; Nooner and Chadwick, 2009] and similar drops are seen over magma chambers during dike propagation events in Iceland [e.g. Einarsson, 1991]. We see no such events in the BPR data suggesting no eruption.

<table>
<thead>
<tr>
<th>Table 1. BPR Deployment Locations</th>
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<tbody>
<tr>
<td>Benchmark</td>
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<tr>
<td>-----------</td>
</tr>
<tr>
<td>BPR08</td>
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<tr>
<td>BPR08</td>
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<tr>
<td>BPR08</td>
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<td>BPR01</td>
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<tr>
<td>BPR06</td>
</tr>
</tbody>
</table>
or dike events occurred from March 2007 until at least November 2010 in the vicinity of the instruments. Over this short-timescale, drift of the gauges and oceanographic signals (other than the tide signal which can be removed) produce negligible change in pressure and thus we can conclude that no rapid elevation changes with amplitudes more than a few centimeter have occurred during the observation interval. The batteries died in BPR01 and BPR06 before we recovered the instruments in October, 2011.

BPR observations at one additional site about 30km to the north (10°4.7282'N, 104°20.2179'W), and three sites at about 27km (9°35.0257'N, 104°15.3048'W), 35km (9°30.3978'N, 104°14.7812'W), and 56km (9°19.0'N, 104°13.0'W) south of the benchmarks also show no evidence for rapid vertical elevation changes that would be indicative of eruptions elsewhere on this ridge segment. These data extend from May 2007 until about mid October 2011.

MPR surveys were carried out during research cruises on board the RV Atlantis in June 2008, December 2009, and October 2011. Before MPR measurements were made in 2008, 10 concrete benchmarks (Figure 3) weighing approximately 91 kg in water were installed on the seafloor. The benchmarks span a range of distances from the AST in order to allow any observed displacement to be compared to deformation models. Benchmark names and locations are shown in Figure 1 and site characteristics are given in Table 2. The benchmarks far from the AST were located on a layer of pelagic sediments that increased in thickness with distance from the AST. The benchmarks legs were pushed completely into the mud (Figure 3a) using the hydraulic manipulator arm of the HOV Alvin during installation. The first two surveys were conducted using the manned DSV Alvin and the final survey was conducted using the ROV Jason. Although it was not our

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**Figure 2.** Filtered BPR data from at the study site showing (a) BPR08 at site EPR07 from March 2007 to April 2008, (b) BPR08 at site EPR09 from July 2008 through December 2009, (c) BPR01 at EPR10 from January 2009 to December 2009, and (d) BPR06 at site EPR01 from January 2009 to August 2010. The BPR pressure data have been converted to depth and tidal frequencies have been filtered out using a low-pass Butterworth filter.
original intent, the use of the HOV Alvin rather than the ROV Jason meant the first two surveys became a test of our MPR instrument limitations and survey techniques, and as such provided us with valuable lessons on survey design and instrument limitations, as described in the next section.

The MPR pressure data were reduced by first removing the effect of tides and then calculating and removing instrument drift. Continuous BPR data collected at the AST (Figure 2, shown with tides removed) were used to correct the MPR data for ocean tides in years 2007 and 2009. The BPRs were no longer recording during the 2011 survey, so the SPOTL tide model was used to make ocean tide corrections [Agnew, 1997]. The BPR data showed good agreement with the SPOTL model in the first two surveys. Following the methods described in Chadwick et al. [2006] and Nooner [2005], the linear drift of the MPR instrument was determined each year from repeated measurements on each benchmark. After correcting for ocean tides and linear instrument drift, the pressure data were then converted to depth using the standard ocean depth formula derived by Fofonoff and Millard [1983]. The resulting uncertainty of the MPR measurements is given by the scatter of repeated measurement on

<table>
<thead>
<tr>
<th>Benchmark</th>
<th>Benchmark Locations</th>
<th>Depth (m)</th>
<th>Number of Measurements Each Year</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Lat. (N)</td>
<td>Lon. (W)</td>
<td>2008</td>
</tr>
<tr>
<td>EPR01</td>
<td>9.85338</td>
<td>104.20995</td>
<td>2787</td>
</tr>
<tr>
<td>EPR02</td>
<td>9.85005</td>
<td>104.22802</td>
<td>2851</td>
</tr>
<tr>
<td>EPR03</td>
<td>9.84732</td>
<td>104.24635</td>
<td>2688</td>
</tr>
<tr>
<td>EPR04</td>
<td>9.84410</td>
<td>104.26388</td>
<td>2611</td>
</tr>
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</tr>
<tr>
<td>EPR06</td>
<td>9.84107</td>
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<td>104.29063</td>
<td>2501</td>
</tr>
<tr>
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<td>2520</td>
</tr>
<tr>
<td>EPR09</td>
<td>9.82970</td>
<td>104.28795</td>
<td>2506</td>
</tr>
<tr>
<td>EPR10</td>
<td>9.84880</td>
<td>104.29211</td>
<td>2509</td>
</tr>
</tbody>
</table>

Table 2. MPR site Characteristics
all the benchmarks for each survey and is 3.5, 2.2, and 0.6 cm in 2008, 2009, and 2011, respectively (Table 3; all uncertainties given are one σ).

Our results in 2008 are 7 times worse than the results from Stenvold et al. [2006] and more than 3–4 times worse than recent results from Axial Seamount [Chadwick et al., 2012, 2006; Nooner and Chadwick, 2009], due to temperature changes in the instrument that occurred during each short HOV Alvin dive and drift correction errors. This is discussed in more detail in a later section. Our results in 2011, however, are as good or better than measurements that have made anywhere else, indicating that high-quality data are achievable at the EPR. In the following three sections, we discuss each of the three surveys in detail.

2.1. Survey One

Although Paroscientific 410K pressure gauges are temperature compensated, rapid changes in temperature have been shown to impact their pressure output value [e.g. Nooner, 2005]—there can be a time lag between the change in temperature of the gauge and the correction made for temperature change. For example, Chadwick et al. [2006] found that it took ~2 h for the pressure gauges to equilibrate once on bottom at Axial Seamount. Previous surveys in the North Sea, however, were able to achieve 0.5 mm measurement repeatability without waiting additional time on bottom [Sasagawa et al., 2003; Stenvold et al., 2006]. In fact, in these North Sea surveys the ROV was usually recovered to deck and the instruments were placed in a bath of temperature-controlled ice water before transiting to the next benchmark. This method was efficient because the water depth ranged only from 80 to 300 m, and launches and recoveries of the ROVs took little time (~10 min). Because of this previous work, we were hopeful that similar results could be obtained in the deeper water over the EPR using the DSV Alvin.

During our survey, the MPR instrument was kept in an ice bath while on deck. It was taken out of the ice bath just prior to launch and was put back into the ice bath immediately following vehicle recovery, following the procedure of Sasagawa et al. [2003] and Stenvold et al. [2006]. Unfortunately, we underestimated (1) the time Alvin would spend in the warm surface waters of the EPR and (2) how rapidly the MPR would warm up, even after spending >12 hours in an ice bath. Each Alvin dive in 2008 took the MPR from an internal temperature of almost 19°C to bottom temperatures as low as 2°C and back during an 8 h period (Figure 4). The short periods of temperatures approaching 2°C that can be seen in Figure 4 result from power being turned off to the MPR instrument during transits. There is a small offset in the location of the temperature sensor and the pressure oscillator within the Paroscientific transducers, resulting in poor temperature corrections during periods of rapid temperature changes (this is discussed in more detail below). By comparison, internal temperature changes during similar surveys in the North Sea were only a few degrees...
Celsius [Stenvold et al., 2006]. We attribute our observed uncertainties of 3.5 cm (T3) to the large and rapid temperature fluctuation the MPR experienced and to poor drift corrections. This level of uncertainty is much higher than we expected based on previous studies. The drift correction uncertainties were high because attempts to complete each dive as a closed loop were not possible given the station spacing and limited bottom time of the DSV Alvin, particularly in the 2008 survey. Most of the nine Alvin dives were shared with another investigator, making it even more difficult to complete a closed loop survey during most dives, and providing another reason that the MPR was periodically powered down. During off axis dives with long transits, we typically ran out of time or battery power in the submersible before being able to complete a loop. Attempts to develop a temperature correction curve did not reduce the repeatability of the overall survey.

2.2. Survey Two
The second survey was carried out in December, 2009. Because of our results from the first survey, we were scheduled to use the ROV Jason for the second survey but logistical problems forced us to use the DSV Alvin again on short notice. We used three strategies to improve our results over the first survey: (1) We carefully planned each dive to optimize the number of repeats during the dive. This allowed us to more accurately estimate a temperature correction curve that we could apply to each dive. (2) We mounted an ice-filled box to the basket of the DSV Alvin to contain the MPR during the dive descent. This isolated the gauges from warming in the surface water prior to descent, thereby minimizing the temperature fluctuations between the ice bath on the deck of the ship and the 2°C bottom waters of the EPR (Figure 4). (3) The last major strategy to improve data quality was that in addition to our normal MPR, attached via cable to the vehicle, we deployed a small battery powered MPR on the last benchmark of each dive. We then started the next dive at that benchmark. This alternate MPR was meant to be the reference gauge that maintained a stable temperature throughout the duration of the survey. This was an instrument that we put together in the week prior to the cruise from spare parts and an unused pressure case that we repurposed. Although it passed the Alvin certification procedure prior to the cruise, the pressure case flooded a few days into the survey and no data were recovered.

The resulting 2009 benchmark height uncertainties are estimated to be 2.2 cm (Table 3). This is better than the previous survey, due primarily to the addition of the ice-filled box on the front basket of Alvin that the MPR was kept in during submersible deployment and descent. We were also able to plan more dives that contained repeated measurements, allowing us to better determine the instrument drift. The single dive repeatability in 2009 was 0.9 cm based on the final dive, which had five repeats, showing that it is possible to carry out a high-quality survey using Alvin with proper station spacing. A linear correction taken from this dive was used to correct the other dives for temperature.

2.3. Survey Three
The ROV Jason was used for the final survey, which took place in October, 2011. In spite of a hurricane, a broken Jason/Medea tether, and a winch failure (all of which consumed considerable ship time and significantly decreased the number of measurements we were able to make) we achieved excellent results (Table 3) with a repeatability of 0.6 cm during two dives. It must be noted that although benchmarks 2, 8, 9, and 10 received only one visit in 2011, benchmark 7 received five visits and all others received two visits each. Uncertainties were calculated based on the repeats of these six stations (Table 2). Temperature records are shown as black circles in Figure 4.

2.4. Temperature Corrections
Pressure is measured within the Paroscientific 410K pressure transducers by utilizing an oscillating quartz crystal beam than is

<table>
<thead>
<tr>
<th>Table 3. Error Budget</th>
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<tr>
<td>Error Source</td>
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<tr>
<td></td>
</tr>
<tr>
<td>Inherent precision of pressure sensor</td>
</tr>
<tr>
<td>Background noise</td>
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<tr>
<td>Calibration uncertainty</td>
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<tr>
<td>Tide correction uncertainty</td>
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<td>Drift correction</td>
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<td>Depth conversion</td>
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<tr>
<td>Rotational uncertainty</td>
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<td>Temperature uncertainty</td>
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<tr>
<td>Root sum of squared errors (expected repeatability)</td>
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<td>Observed repeatability</td>
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mechanically connected to a C-shaped Bourdon tube (http://www.paroscientific.com). Changes in fluid pressures within the Bourdon tube and temperature changes of the quartz beam result in changes in the oscillation frequency of the quartz beam. Since we are only interested in measuring changes in pressure, a temperature correction must be made. To make this correction, temperature within the transducers is measured using a quartz crystal torsional tuning fork temperature sensor (http://www.paroscientific.com). Since the quartz beam and the torsional tuning fork cannot be located in the exact same place, the temperature being measured is at a slightly different location than the quartz beam.

There are two sources of error resulting from the corrections described above: (1) Due to the small offset between temperature sensor and pressure sensor, temperature gradients can cause the temperature and pressure sensors to experience different temperatures. Due to the thermal mass of the instrument, large changes in temperature, like what we observe from the ocean surface to the seafloor, take some time to fully equilibrate. The larger the temperature change is, the longer the equilibration time. While the instrument is equilibrating, there is a changing temperature gradient that can affect the pressure output. (2) Changes in gauge temperature most likely result in a change in the shape of the Bourdon tube itself due to thermal expansion or contraction. This results in small differences in the apparent pressure.

The temperature records for the three surveys shown in Figure 4 show a similar broad-scale trend. In general, the MPR instrument begins a dive cool (since it was kept in a bucket of ice water prior to the dives), then goes through a period of rapid warm-up once contact is made with the ocean surface waters. This is followed by a period of rapid cool down as the outside water temperature decreases with depth, and finally a period of relatively stable temperatures. The Jason dives show a much more repeatable temperature profile than the Alvin dives do, which show particularly large variability in 2008 before we used an ice-box on the science basket to carry down the MPR. The temperature of the instrument during the Jason dives stabilizes after about 3 h (it takes about 1.5 h to reach the bottom) with very few fluctuations for the rest of the dive (only the first 10 h of each dive are shown). The temperature of the instrument during the Alvin dives is similar, but with much more variability. Part of this variability in gauge temperature is due to MPR power being turned off to conserve power during Alvin transits. In retrospect, this was not good practice since loss of power to the MPR allowed the instrument to cool down by almost 2°C. Power was restored to make the subsequent measurement, causing the MPR electronics to heat the interior of the instrument and inducing a temperature gradient in the pressure transducers. Aside from a single dive, the MPR power was kept on for the duration of each dive during the 2009 survey.

The Paroscientific 410K pressure transducers we used in these surveys have temperature corrections factors that range from 20 to 120 cm/°C at the survey depths. The following exercise illustrates the magnitude of the errors we might expect from turning the MPR power off during transits, as was done in much of the 2008 survey. The pressure transducers are enclosed inside a cylindrical aluminum pressure case (Figure 3) that is about 16 cm in diameter. If we assume that the aluminum walls are at 2°C and the MPR electronics are at 4°C, then a rough estimate of the temperature gradient from the center to the walls of the MPR is 0.25°C/cm. Therefore, if the torsional tuning fork is 0.1 cm from the quartz beam (this is an estimate only), we would expect that the temperature difference between the two locations would be about 0.025°C. The resulting temperature correction translates to depths errors ranging from 0.5 to 3 cm, depending on the gauge.

The very small temperature uncertainty for 2011 (Table 3) results from the fact that the majority of the measurements during the two ~24 h Jason dives were made significantly after the 3 h equilibration time after the dives began and there were only two dives. In addition to the temperature correction errors discussed above, the 2008 and 2009 surveys were both carried out over the course of nine Alvin dives, during which at least 1/4 of the measurements were made before the gauges had equilibrated.

3. Discussion

Figure 5 shows the benchmark heights each year relative to the reference benchmark EPR01. The uncertainty in benchmark height changes from 2008 to 2009 is 4.1 cm—the square root of the sum of the squared uncertainties for each year. Within this uncertainty, we cannot be confident of any benchmark height changes during that time. However, this null result puts an upper bound of about 8.1 cm on the
maximum deformation that could have occurred at 9°50′ N during that time window. From December 2009 to October 2011, we observe up to 12 cm of uplift on benchmark EPR10 (Figure 5). This exceeds the combined 2.3 cm uncertainty for the 2008–2009 time span by a factor of 5, giving us confidence in our observations. Uplift generally decreases with distance from EPR10 (Figures 5 and 6), suggesting that the deformation results from volcanic inflation at the AST, however, benchmarks EPR02 and EPR03 show anomalous uplift compared to this trend. This may be caused by tectonics (i.e., faulting) but is more likely due to either measurement errors in 2008 or benchmark instabilities.

We fit the observed deformation from 2009 to 2011 using a point source inflation model [Mogi, 1958] with a Poisson’s ratio of $\nu = 0.296$, which we calculated from $V_p/V_s$ ratios that have been measured for depths of 1–3 km at this location [e.g., Vera et al., 1990; Waldhauser and Tolstoy, 2011]. Decreasing Poisson’s ratio decreases the volume change of the optimum inflation source, while the depth remains the same. If we remove the results from benchmarks EPR02 and EPR03, the deformation data can be fit with a point source located on the ridge axis 407 m northeast of our benchmark EPR10 (best fit is at 9°51.166′ N and 104°17.540′ W; Figures 1 and 6). This location roughly coincides with a fourth-order discontinuity in the ridge morphology and the boundary between two seismically imaged shallow axial melt lens segments from Carbotte et al. [2013]. The best fitting point source depth is 2.7 km, whereas the magma lenses from Carbotte et al. [2013] lie at about 1.5 km depth at this location. Interestingly, Marjanovic [2013] and Marjanovic et al. [2013] found a region of enhanced melt underlying the melt lens over most of this segment of the EPR, except in the location of our best fitting inflation source. They showed that melt present in the midcrust underlying this region of the EPR from 9°50′ to 9°52′ N is
significantly lower than adjacent segments. Bathymetric and geochemical data [Goss et al., 2010; Soule et al., 2007; White et al., 2006] indicate that the source of the eruption was 9°48′ - 9°52′ N. Marjanovic et al. [2013] suggested that the 2005–2006 eruption was sourced initially from the midcrustal melt body 9°50′–9°52′ N, depleting the midcrust of melt in this region. Our geodetic results suggest that inflation could be occurring at a depth of about 2.7 km underlying the 9°51′ N region, consistent with a model of reinflation for a midcrustal magma region that drained due to the eruption. It is important to point out that a similar deformation pattern could be obtained for a shallower axisymmetric oblate spheroidal inflation source (a penny-shaped crack) if the ratio of radius to depth is greater than about 0.8 [Fialko et al., 2001]. This would mean that a 1.5 km deep penny-shaped magma chamber would have to be around 3 km wide; however, this analysis does not reflect the along-axis dimension of the observed magma lens [Carbotte et al., 2013].

Our results suggest that magmatic reinflation is occurring at 9°50′ N at a rate of 7 cm/yr, equivalent to a point source volume change of $2.3 \times 10^6$ m$^3$/yr. This is about 1/13 the rate of long-term inflation observed at Axial Seamount [Chadwick et al., 2006; Nooner and Chadwick, 2009]. Axial Seamount is a large volcanic seamount associated with the Cobb hotspot, so its magma supply is likely enhanced relative to a normal fast spreading mid-ocean ridge segment like the EPR at 9°50′ N. Rapid reinflation periods following both the 1998 and 2011 eruptions of Axial Seamount have been attributed to viscoelastic relaxation and/or poroelastic behavior of the mush region underlying the magma chamber [Chadwick et al., 2013; Nooner and Chadwick, 2009]. This period of inflation was very brief, lasting only a few months. Therefore, we believe that since the EPR eruption occurred 3–4 years prior to the observed uplift shown here, this phase of inflation was already past. If the observed inflation rate at the EPR represents a steady linear long-term inflation, we expect the magma chamber to be fully reinflated and pruned for an eruption by 2026, after an approximately 20 year interval. We obtain this result by using $47 \times 10^6$ m$^3$ as the total volume of magma evacuated from the magma chamber during the 2005/6 eruption from Soule et al. [2007], and assuming a continuous inflation rate of $2.3 \times 10^6$ m$^3$/yr. This estimate is in general agreement with the known recurrence interval of 13 years between the 1992 and the 2005/2006 eruptions. The volume of the 2005/2006 eruption is only 25–33% of the volume of the two eruptions observed at Axial Seamount. This EPR site has a much smaller magma chamber than Axial Seamount as inferred from seismic studies [e.g., Detrick et al., 1987; Soule et al., 2007; West et al., 2001].

The only other constraints on magmatic inflation along spreading segments come from recent observations at two subaerial regions. The estimated inflation rate at 9°50′ N on the EPR is about a factor of 3 lower than the average 20 cm/yr inflation rate at the center of the Krafla spreading segment in Iceland over a 10 year episode of activity that began in 1975 [e.g., Tryggvason, 1984]. However, there were 20 dike intrusion and/or extrusion events in that episode over that period on this slow spreading segment. Also, the time interval since the last episode on the Krafla segment occurred 250 years before the 1975 sequence [e.g., Einarssson, 1991]. The only other constraint on magmatic reinflation on a spreading center comes from an episode of 14 dike intrusion events on the Dabbahu spreading segment in the Afar of Ethiopia that began in 2005. Six years of InSAR observations there show about 20 cm/yr of uplift of the segment center which is comparable to the rate seen for Krafla [e.g., Wright et al., 2012].

Our experiences at the EPR suggest that it is possible to carry out an MPR survey using the DSV Alvin; however, certain special precautions must be made. (1) Stowing the MPR inside a sealed ice-box during vehicle launch greatly reduces the temperature changes that the MPR experiences. This leads directly to lower uncertainties during the short Alvin dives because it reduces the temperature correction lag and shortens the time for the MPR to equilibrate with the water temperature near the seafloor. (2) Stations should be arranged to allow at least one repeat measurement during each dive. This will allow a combined temperature and drift correction to be made for each dive separately, as required by our observation that this correction varies from dive to dive. (3) Finally, an autonomous pressure gauge package should be used as an additional MPR and left on the seafloor the duration of the survey to add redundancy. This gauge would then be left on the final benchmark at the end of each dive and picked up at the beginning of the subsequent dive. Data from this gauge would then be the same as a typical ROV-based MPR survey.

4. Conclusions

In spite of challenging circumstances during each cruise, resulting in data uncertainties much greater than we expected, we have strong evidence that up to 12 cm of inflation (7 cm/yr) occurred at 9°50′ N on the
EPR between December 2009 and October 2011. The focus of this inflation is 407 m NNE of EPR10, on the ridge axis at a depth of 2.7 km. This location and depth coincide with the location of a magmatically depleted midcrustal region that Marjanovic et al. [2013] observed following the eruption. The magma supply rate required is about half what has been observed as the "steady state" magma supply rate at Axial Seamount on the Juan de Fuca Ridge, suggesting a less vigorous magma supply underlies this part of the EPR.

The eruption occurred in 2005/2006, but our first survey was not conducted until 2008, more than 2 years after the eruption. We hypothesize that the observed inflation rate of 7 cm/yr represents the "steady state" supply from a deeper melt reservoir. This hypothesis can be tested via additional pressure measurements on the geodetic network that we have established.

The BPR data suggest that no new eruptions have occurred in the vicinity of the 2005 eruption between March 2007 and October 2011, nor have eruptions or dike intrusions occurred elsewhere along this ridge segment within the vicinity of one BPR 30 km to the north of the eruption site, or near three BPRs 27 to 50 km south of the site.

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