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Periodic solutions of arbitrary length in a simple integer iteration

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PERIODIC SOLUTIONS OF ARBITRARY LENGTH IN A SIMPLE INTEGER ITERATION

DEAN CLARK

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We prove that all solutions to the nonlinear second-order difference equation in integers *y*_{*n*+1} = $[ay_n] - y_{n-1}$, {*a* ∈ ℝ : |*a*| < 2, *a* ≠ 0,±1}, *y*₀, *y*₁ ∈ ℤ, are periodic. The first-order system representation of this equation is shown to have self-similar and chaotic solutions in the integer plane.

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1. Introduction

We study the nonlinear second-order difference equation in integers

$$
y_{n+1} = [ay_n] - y_{n-1}, \quad \{a \in \mathbb{R} : |a| < 2, \, a \neq 0, \pm 1\}, \, y_0, y_1 \in \mathbb{Z}, \tag{1.1}
$$

where $\lceil x \rceil$ denotes the smallest integer not smaller than x (the *ceiling* function). The reader is already familiar with the linear cases $a = 0, \pm 1$, therefore we do not consider these values in this paper. Besides the natural generalization to discrete space, there are at least three reasons why [\(1.1\)](#page-2-0) is interesting.

First, when $a = 3/2$, [\(1.1\)](#page-2-0) becomes

$$
y_{n+1} = \begin{cases} \frac{3y_n + 1}{2} - y_{n-1} & \text{if } y_n \text{ is odd} \\ 3\left(\frac{y_n}{2}\right) - y_{n-1} & \text{if } y_n \text{ is even,} \end{cases}
$$
(1.2)

a second-order variant of the notorious "3*x* + 1 iteration." So far as we know, the ultimate convergence to 1 of the $3x + 1$ iterates remains an unproven conjecture. In contrast, we will prove an ultimate recurrence property for [\(1.1\)](#page-2-0) for all initial states $y_0, y_1 \in \mathbb{Z}$ and parameter values [−]² *<a<* 2. It is the initial state that is always recurrent. Moreover, solutions to [\(1.1\)](#page-2-0) can exhibit periods of arbitrary length (Theorems [2.2,](#page-3-0) [3.2,](#page-6-0) below).

Hindawi Publishing Corporation Advances in Difference Equations Volume 2006, Article ID 35847, Pages 1[–9](#page-10-0) DOI [10.1155/ADE/2006/35847](http://dx.doi.org/10.1155/S1687183905505053)

Second, the method used to establish the periodicity of all solutions to [\(1.1\)](#page-2-0) seems novel, elegant, and of potentially wider applicability. Without this method, we could not prove that solutions of the special case [\(1.2\)](#page-2-1), above, were even bounded [\[1](#page-10-1)].

Third, [\(1.1\)](#page-2-0) is converted to a first-order system in two variables using the mapping $T(x, y) = (y, \lceil ay \rceil - x)$. The simplicity of *T* gives no hint of the startling complexity shown by scatter plots of some of the solutions. See Figures [4.1](#page-8-0)[–4.4,](#page-10-2) below.

2. Qualitative behavior of solutions

Henceforth, all pairs (x, y) denote points in the integer plane \mathbb{Z}^2 , and the real parameter *^a* satisfies [|]*a*[|] *<* 2. We obtain the aforementioned first-order system by letting

$$
T(x, y) = (y, \lceil ay \rceil - x), \qquad X_n = (x_n, y_n) = T^n(x_0, y_0), \quad n = 0, 1, 2, \dots \tag{2.1}
$$

Remark 2.1. The sequence (y_n) appearing as the second coordinate in each term of (X_n) is the same sequence generated by [\(1.1\)](#page-2-0) when $x_0 = \lceil ay_0 \rceil - y_1$.

A first glimpse of the rotational motion of solutions is obtained from the powers of $A = \begin{pmatrix} 0 & 1 \\ -1 & a \end{pmatrix}$, the matrix underlying *T* without the ceiling function. Because $|a| < 2$, *A* has complex eigenvalues. After diagonalizing *A*, we have

$$
A^{n} = \begin{pmatrix} \cos(n\theta) - \frac{a\sin(n\theta)}{\sqrt{4 - a^{2}}} & \frac{2\sin(n\theta)}{\sqrt{4 - a^{2}}} \\ -\frac{2\sin(n\theta)}{\sqrt{4 - a^{2}}} & \cos(n\theta) + \frac{a\sin(n\theta)}{\sqrt{4 - a^{2}}} \end{pmatrix}, \quad \theta = \arccos\left(\frac{a}{2}\right). \tag{2.2}
$$

The significance of θ for the nonlinear equation [\(1.1\)](#page-2-0) will become apparent later.

The identity $x^2 + y^2 - axy = y^2 + (ay - x)^2 - ay(ay - x)$ supplies a family of *invariant ellipses* $E(x, y) = x^2 + y^2 - axy$ for the *linear* equation. [Figure 2.1](#page-4-0) shows the ellipse x^2 + $y^2 - (1/2)x y$ determined by $a = 1/2$ and $(x_0, y_0) = (0, 32)$, as well as the first six iterates of *^T* acting on (*x*0, *^y*0) ⁼ (0,32) : (0,32), (32,16), (16,−24), (−24,−28), (−28,10), and (10,33). All these points lie on the ellipse $x^2 + y^2 - (1/2)xy = 1024$ because the ceiling function is inactive. The first odd y_n requiring use of the ceiling is $y_5 = 33$, and we expect that $T(10,33) = (33,7)$ does not lie on this ellipse. Indeed, it does not: $33^2 + 7^2 - 1/2$. $33 \cdot 7 = 1022.5$.

The clockwise motion in \mathbb{Z}^2 of the iterates of *T* is further clarified by the vector field drawn in [Figure 2.2](#page-4-1) with $a = 1/3$. Each directed segment at (x, y) has the form $T(x, y)$ – (x, y) and thus points toward $T(x, y)$. This is seen clearly in [Figure 2.2](#page-4-1) for the orbit initiated at (1,0). The vector field is divided into four quadrants with boundaries $y = x$ and the step locus $y = [ay] - x$. Each linear segment of this locus includes the upper lefthand endpoint and excludes the lower right-hand endpoint. The quadrants discriminate whether direction vectors have $\Delta x > 0 \leq 0$ or $\Delta y > 0 \leq 0$. In general, the clockwise rotation and roughly elliptic orbits are easily confirmed for all parameter values [−]² *<a<* ² and initial conditions $(x_0, y_0) \neq T(x_0, y_0)$.

Figure 2.1. $a = 1/2$, $X_0 = (0, 32)$.

Figure 2.2. $a = 1/3$, $X_0 = (1,0)$.

THEOREM 2.2. *For nonzero rational* $a = p/q$ *where* p *and* q *have no common factors, the number of distinct terms of a solution* (*yn*) *can be made arbitrarily large depending on the initial conditions.*

Proof. Imitating the example of [Figure 2.1,](#page-4-0) we choose initial conditions designed to deactivate the ceiling function for a finite number of terms, thereby turning the nonlinear equation [\(1.1\)](#page-2-0) into a linear equation. Take $y_0 = q^m$, $y_1 = pq^{m-1}$ with arbitrarily large

positive *m*. As before, *A* denotes the matrix underlying the linear system. By induction,

$$
A^{n} = \begin{pmatrix} \frac{-f_{n-2}}{q^{n-2}} & \frac{f_{n-1}}{q^{n-1}}\\ \frac{-f_{n-1}}{q^{n-1}} & \frac{f_n}{q^n} \end{pmatrix},
$$
(2.3)

where $f_{-1} = 0$, $f_0 = 1$, and $f_n(p,q) = \sum_{k=0}^{\lfloor n/2 \rfloor} {n-k \choose k} (-1)^k p^{n-2k} q^{2k}$ for $n > 0$. A consequence of taking *p* and *q* relatively prime is that *q* never divides f_n for $n \ge 0$; the coefficient of p^n in f_n is always 1. Repeated application of *A* to the initial vector $(0, q^m)$ gives the general form of the first *m* + 1 iterates: $y_n = f_n q^{m-n}$. These are all distinct because the highest power of *a* that divides each one is different. □ power of *q* that divides each one is different.

In contrast, the following example shows that when *a* is *irrational*, it does not follow that there are arbitrarily many distinct iterates simply because an initial condition is arbitrarily large (however, see [Figure 4.1](#page-8-0) below).

Example 2.3. Let $a = (\sqrt{5} - 1)/2 = 0.6180339...$, that is, $\theta = 2\pi/5$ in [\(2.2\)](#page-3-1), above. Let $y_0 = 1$, $y_1 = 10^n$ for $n \ge 0$. With this form of initial condition, all solutions have period 5. Solutions are shown, below, for $n = 0, 1, 2, 3$, and 6.

The curious relation between y_2 and y_3 should be noted: sometimes $y_3 = -y_2$ and sometimes $y_3 = -y_2 - 1$. It is easy to see that $y_2 = \lfloor 10^n a \rfloor$, where $\lfloor x \rfloor$ denotes the greatest integer not greater than *x*. With a little more work, using the fact that $a^2 + a - 1 = 0$, we can prove that $y_3 = -y_2$ if and only if $1 - a < \lceil 10^n a \rceil - 10^n a$; otherwise, $y_3 = -y_2 - 1$.

3. Periodicity of solutions

An *involution* is a map *V* such that the square of *V* is the identity, that is, $V^2 = V \cdot V = I$ [\[2\]](#page-10-3). The following lemma provides basic machinery for proving that all solutions of [\(1.1\)](#page-2-0) are periodic.

LEMMA 3.1. *Let T be defined as in* [\(2.1\)](#page-3-2), above, and S(*x*, *y*) ≡ $T^{-1}(x, y) = (\lceil ax \rceil - y, x)$. The *involution* $V(x, y) = (y, x)$ *satisfies* $VT = SV$ *and* $TV = VS$ *. It follows that the mappings VT, VS, TV, and SV are involutions.*

Proof. We have

$$
V(T(x, y)) = V(y, [ay] - x) = ([ay] - x, y) = S(y, x) = S(V(x, y)).
$$
 (3.1)

Multiplying $VT = SV$ on the left and right by *V* yields $TV = VS$, which is used to prove that $VTVT = VVST = I$. Thus, VT is an involution and similarly so are *VS*, TV , \Box and *SV*.

We call a solution of [\(1.1\)](#page-2-0) *invariant under V* if the point set $O = \{X_n \in \mathbb{Z}^2 : X_n =$ *Tn*(*X*₀), *n* = 0,1,2,...} satisfies *V*(*O*) = *O*. Geometrically, for *V*(*x*, *y*) = (*y*,*x*), this means that the plot of iterates is symmetric about the 45◦ line, for example, see [Figure 2.2,](#page-4-1) above. For *rational a*, numerical experiments have shown that this invariance is so prevalent, we conjecture it occurs with probability 1. See the corollary to [Theorem 3.2,](#page-6-0) below. For solutions invariant under *V*, the lemma establishes periodicity at once:

$$
X_0 = TVTV(X_0) = TVT(X_k) = TV(X_{k+1}) = T(X_m) = X_{m+1}.
$$
\n(3.2)

The general periodicity result follows, with [Figure 3.1,](#page-7-0) below, providing concrete support to the proof.

THEOREM 3.2. *For* $a \in \mathbb{R}$, $|a| < 2$, *all solutions of* [\(1.1\)](#page-2-0) *are periodic.*

Proof. Let a solution to [\(1.1\)](#page-2-0) begin y_0, y_1, \ldots Citing [Remark 2.1,](#page-3-3) take $X_0 = (\lceil ay_0 \rceil - y_1, y_0)$ in \mathbb{Z}^2 . The mappings *T*, *S*, and *V* are defined as in [\(2.1\)](#page-3-2) and [Lemma 3.1,](#page-5-0) respectively. Set $Y_0 = V(X_0)$ and $Y_{k(n)} = V(X_n)$ for $n = 1, 2, \ldots$ The value of *k* is determined by use of the lemma in [\(3.3\)](#page-6-1), below: *k* is the number of times *S* must be applied to the point $V(X_n)$ so that the points Y_k , Y_{k-1} , Y_{k-2} ,... rotate (counterclockwise) back to Y_0 . See [Figure 3.1.](#page-7-0)

$$
Y_{k-1} = S(Y_k) = SV(X_n) = VT(X_n) = V(X_{n+1}),
$$

\n
$$
Y_{k+1} = T(Y_k) = TV(X_n) = VS(X_n) = V(X_{n-1}).
$$
\n(3.3)

Again, by the lemma, [\(3.3\)](#page-6-1) implies that *n* applications of *T* to $V(X_n)$ move Y_k, Y_{k+1} , Y_{k+2} ,... (clockwise) to $Y_{k+n} = V(X_0) = Y_0$. Thus the sequence (Y_k) is periodic. By definition, (X_n) and (Y_k) are in one-to-one correspondence by way of reflection across the 45[°] line. Thus, (X_n) is periodic. In particular, [\(3.3\)](#page-6-1) implies $Y_0 = V(X_0) = V(X_{n+k})$; hence, $X_0 = X_{n+k}$. In accordance with [Remark 2.1,](#page-3-3) all solutions to [\(1.1\)](#page-2-0) are periodic.

In [Figure 3.1](#page-7-0) the points of (X_n)

$$
(2,-3), (-3,-6), (-6,-5), (-5,-1), (-1,4), (4,7), (7,6), (6,2), (2,-3) \tag{3.4}
$$

are denoted by black circles. The points of (Y_k) , which are read from right to left in [\(3.4\)](#page-6-2) with *V* applied to each pair, are denoted by open circles in [Figure 3.1.](#page-7-0) The following corollary deals with the special case where the initial pair lies on the 45◦ line.

COROLLARY 3.3. *For* $a \in \mathbb{R}$, $|a| < 2$, and $X_0 = (y_0, y_0)$ all solutions of [\(1.1\)](#page-2-0) are invariant *under V.*

Figure 3.1. The subscripts of *Xn* and its image under *V* always sum to 8.

Figure 3.2. *^X*⁰ ⁼ (20,−30), *^a* ⁼ ⁷*/*⁵ ⁼ ¹*.*4 (black squares); *^a* ⁼ [√] 2 (open circles).

Proof. Solutions are periodic by [Theorem 3.2.](#page-6-0) Suppose that, for a given *a* and $X_0 = (y_0,$ y_0), the resulting solution has smallest period *N*, so that $X_0 = X_N$. Since $V(X) = X$ on the 45◦ line, the lemma yields

$$
X_1 = T(X_0) = TV(X_0) = VS(X_N) = V(X_{N-1}).
$$
\n(3.5)

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Figure 4.1. $a = (\sqrt{5} - 1)/2$, $\theta = 2\pi/5$.

Next,

$$
X_2 = T(X_1) = TV(X_{N-1}) = VS(X_{N-1}) = V(X_{N-2}).
$$
\n(3.6)

Continuing in this way, $X_k = V(X_{N-k})$ for $k = 0, 1, 2, ..., N$. \Box

By re-indexing, it is clear that if any iterate touches the 45◦ line, the entire trajectory becomes symmetric about this line. Perhaps this explains why there are so many invariant solutions when *a* is rational. The rotation angle $\theta = \arccos(a/2)$ ([\(2.2\)](#page-3-1), above) is never a rational multiple of π for nontrivial rational *a*, that is, $a \neq 0, \pm 1$ [\[3\]](#page-10-4), indicating a large number of iterations relative to the size of *X*0. The more densely packed with points a trajectory is, the greater the likelihood that one of them is located on the 45◦ line. In any event, the small-period non-invariant solution with rational *^a* ⁼ ⁷*/*5 shown in [Figure 3.1,](#page-7-0) above, was found by observing that $7/5 = 1.4$ is a fair approximation of $\sqrt{2} = 2\cos(2\pi/8)$ for a small initial point; hence, the period-8 solution [\(3.4\)](#page-6-2). Predictably, with the same $a = 7/5$ and larger $X_0 = (20, -30)$, we get the period-79, *V*-invariant solution shown in [Figure 3.2.](#page-7-1) In this solution, indicated by the dark squares, $X_{59} = (60, 60)$. By the corollary, above, the entire orbit must line up with itself when reflected about the 45◦ line. Maintaining *^X*⁰ ⁼ (20,−30) and changing *^a* to [√] 2 gives back a period-8 non-invariant solution shown by the open circles in [Figure 3.2.](#page-7-1)

4. Self-similarity and chaos

The presence of symmetry in a brute iteration, perhaps just by accident of sheer numbers, is striking. Still more improbable is that, as the initial condition becomes larger, such

Figure 4.2. $a = (1 - \sqrt{5})/2$, $\theta = 3\pi/5$.

Figure 4.3. *^a* ⁼ 2cos(2*π/*7), *^θ* ⁼ ²*π/*7.

a process can generate images with self-similar complexity as it winds blindly around the plane. See Figures [4.1](#page-8-0) and [4.2.](#page-9-0) Each of Figures [4.1](#page-8-0)[–4.4](#page-10-2) shows a different choice of the parameter *a* and several orbits for each choice. Even chaos is possible for specific initial conditions when $a = 2\cos(\theta)$ and θ is a rational multiple of π. Such solutions give rise to entirely unexpected structures as the initial point gets larger. For instance, [Figure 4.4](#page-10-2) shows just four orbits, with bizarre excrescences forming a single outermost orbit. Where it seems incontrovertible that the fractal stars in [Figure 4.1](#page-8-0) will continue to develop their repetitive complexity, there is no telling what may emerge from the vaguely

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Figure 4.4. *^a* ⁼ 2cos(2*π/*9), *^θ* ⁼ ²*π/*9.

bio-reproductive shapes in [Figure 4.3](#page-9-1) as we zoom out. Evidently, only the distance from the origin of a properly chosen initial pair limits the complexity of these images.

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